# The Crystal Structure of Anhydrous Stannous Phosphate, $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ 

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#### Abstract

Anhydrous stannous phosphate, $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$, crystallizes in the monoclinic space group $P 2_{1} / c$ with $Z=4$. The unit-cell parameters at $25^{\circ} \mathrm{C}$ are $a=11.092$ (8), $b=4.830(4), c=16.401$ (10) $\AA$ and $\beta=94.28(5)^{\circ}$. The structure was solved by the heavy-atom method and refined by full-matrix least-squares techniques to $R_{w}(F)$ $=0.034$ and $R(F)=0.047$ with 1813 reflections. The structure consists of alternating layers of $\mathrm{Sn}^{\prime \prime}$ and $\mathrm{PO}_{4}$ ions parallel to the $a c$ plane. Two open channels parallel to $[010]$ are formed by $\mathrm{Sn}^{11}$ ions arranged in a helical fashion. Each $\mathrm{Sn}^{11}$ ion is at the apex of a trigonal pyramid with the three nearest O atoms, each from a different $\mathrm{PO}_{4}$ group, forming the base. In one case, two $\mathrm{Sn}^{11}$ ions enter into a dimeric configuration by sharing an $\mathrm{O} \cdots \mathrm{O}$ edge of the pyramid. The corresponding $\mathrm{Sn}-\mathrm{O}$ distances are $2 \cdot 230$ ( 6 ) and $2 \cdot 317$ ( 6 ) $\AA$ : the $\mathrm{O}-\mathrm{Sn}-\mathrm{O}$ angle of $69.6(2)^{\circ}$ is unusually small.


## Introduction

Anhydrous stannous phosphate, $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$, was first reported as a product of the reaction of stannous chloride with hydroxyapatite (Collins, Nebergal \& Langer, 1961; Collins, 1962). The compound has also been prepared by heating stannous oxide with either phosphoric acid or phosphorus pentoxide (Jordan, Gerrity \& Hanusa, 1976).

A hydrated form, $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2} \cdot 3 \mathrm{H}_{2} \mathrm{O}$, has been reported (Klement \& Haselbeck, 1963). The powder diagram of this hydrated form is distinctly different from the powder diagram of $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{3}$ (Collins, 1962).

This structure determination is part of a study in which we seek crystal-chemical information about products that may form in reactions between stannous compounds and tooth enamel.

## Experimental

The crystal used in this study was obtained by heating a mixture of $\mathrm{Sn}_{3} \mathrm{~F}_{3} \mathrm{PO}_{4}$ and $\mathrm{Sn}_{2} \mathrm{FPO}_{4}$ that had been made by mixing $1 M$ solutions of $\mathrm{SnF}_{2}$ and $\mathrm{H}_{3} \mathrm{PO}_{4}$. The crystal was approximately rectangular with dimensions $0.048 \times 0.042 \times 0.030 \mathrm{~mm}$. Because of the rather irregular shape of this crystal, realistic absorption corrections could not be applied. Estimated values for the maximum and minimum transmission factors ( $\mu=$ $88.9 \mathrm{~cm}^{-1}$ ) are 0.82 and 0.63 respectively. The sys-

[^0]tematic absence of reflections $h 0 l$ with $l=2 n+1$ and $0 k 0$ with $k=2 n+1$ showed the space group to be $P 2_{1} / c$. The crystal data at room temperature are: $a=$ 11.092 (8), $b=4.830(4), c=16.401$ (10) $\AA, \beta=$ $94.28(5)^{\circ}, D_{c}=4.138 \mathrm{~g} \mathrm{~cm}^{-3}, V=876.2 \AA^{3}, \mathrm{FW}$ $546 \cdot 0, F(000)=976, Z=4$.

The cell dimensions and standard deviations were obtained from least-squares refinements of $202 \theta$ values obtained by automatically centering reflections on a four-circle diffractometer equipped with a graphite monochromator. Mo $K a_{1}(\lambda=0.70930 \AA)$ radiation was used. The diffractometer-controlling program was written by Lenhert (1975).

X-ray intensity data were collected with $\theta-2 \theta$ scans. The $2 \theta$ scan rate was $1^{\circ} \mathrm{min}^{-1}$ and backgrounds were counted for 20 s at both ends of the scan range. All reflections in the hemisphere ( $\pm h k \pm l$ ) up to $2 \theta=80^{\circ}$ were measured, thus giving at least two measurements for each unique reflection. Three standard reflections were measured at intervals of 50 as a check of variation in intensity with time for scaling purposes. The data were averaged into a unique set of 1985 reflections, of which 1813 had observable intensity $|I \geq 2 \sigma(I)|$ and were used in the least-squares refinements. Equivalent reflections agreed within $10 \%$ on the average (based on $F^{2}$ ).

## Structure determination and refinement

The structure was solved by the heavy-atom method. The positions of two Sn atoms were readily deduced from a sharpened, three-dimensional Patterson synthesis. Although the third Sn atom position could also be obtained from the Patterson map, it was actually taken from a Fourier synthesis phased with two Sn

## Table 1. Fractional coordinates and thermal parameters

The figures in parentheses are the estimated standard deviations. The positional parameters are multiplied by $10^{5}$ for the Sn atoms and by $10^{4}$ for the remaining atoms. The thermal parameters are of the form $\exp \left[-2 \pi^{2}\left(U_{11} h^{2} a^{* 2}+U_{22} k^{2} b^{* 2}+U_{33} l^{2} c^{* 2}+2 U_{12} h k a^{*} b^{*}+\right.\right.$ $\left.\left.2 U_{13} h l a^{*} c^{*}+2 U_{23} k l b^{*} c^{*}\right)\right]$, and are multiplied by $10^{4}$ for the Sn and P atoms and by $10^{3}$ for the O atoms.

|  | $x$ | $y$ | $z$ | $U_{11}$ | $U_{22}$ | $U_{33}$ | $U_{12}$ | $U_{13}$ | $U_{23}$ |
| :--- | ---: | :---: | ---: | :---: | :---: | :---: | ---: | ---: | ---: |
| $\mathrm{Sn}(1)$ | $80055(6)$ | $25527(16)$ | $28333(4)$ | $108(3)$ | $137(3)$ | $173(3)$ | $-10(3)$ | $6(3)$ | $-17(3)$ |
| $\mathrm{Sn}(2)$ | $47898(7)$ | $22053(16)$ | $13615(4)$ | $183(4)$ | $174(4)$ | $173(3)$ | $-3(4)$ | $4(3)$ | $-2(3)$ |
| $\mathrm{Sn}(3)$ | $14388(7)$ | $37016(14)$ | $46117(5)$ | $172(3)$ | $141(3)$ | $165(3)$ | $2(4)$ | $19(2)$ | $-20(4)$ |
| $\mathrm{P}(1)$ | $239(2)$ | $2891(5)$ | $1480(1)$ | $132(11)$ | $67(12)$ | $88(10)$ | $-5(10)$ | $2(9)$ | $-17(11)$ |
| $\mathrm{P}(2)$ | $6485(2)$ | $1949(5)$ | $4742(2)$ | $108(12)$ | $94(13)$ | $217(14)$ | $11(11)$ | $12(11)$ | $-25(11)$ |
| $\mathrm{O}(1)$ | $1295(6)$ | $1610(12)$ | $1979(4)$ | $16(4)$ | $7(3)$ | $19(3)$ | $0(3)$ | $-3(3)$ | $-2(3)$ |
| $\mathrm{O}(12)$ | $156(6)$ | $1778(13)$ | $592(3)$ | $14(3)$ | $9(4)$ | $13(3)$ | $3(3)$ | $5(3)$ | $-2(3)$ |
| $\mathrm{O}(13)$ | $-995(6)$ | $2146(14)$ | $1788(4)$ | $17(3)$ | $12(4)$ | $15(3)$ | $-5(3)$ | $8(3)$ | $-1(3)$ |
| $\mathrm{O}(14)$ | $442(6)$ | $6033(14)$ | $1481(4)$ | $17(4)$ | $9(3)$ | $14(3)$ | $-1(3)$ | $-3(3)$ | $-2(3)$ |
| $\mathrm{O}(21)$ | $7753(6)$ | $3113(14)$ | $4610(4)$ | $13(4)$ | $21(4)$ | $36(4)$ | $-6(3)$ | $5(3)$ | $-11(4)$ |
| $\mathrm{O}(22)$ | $5524(6)$ | $3260(13)$ | $4148(4)$ | $17(4)$ | $11(4)$ | $25(4)$ | $-2(3)$ | $-5(3)$ | $-4(3)$ |
| $\mathrm{O}(23)$ | $6239(6)$ | $2374(16)$ | $5636(4)$ | $15(3)$ | $26(4)$ | $24(4)$ | $2(4)$ | $2(3)$ | $-9(4)$ |
| $\mathrm{O}(24)$ | $6495(6)$ | $-1194(14)$ | $4539(4)$ | $9(3)$ | $11(3)$ | $35(4)$ | $3(3)$ | $-4(3)$ | $0(4)$ |

atoms. The remaining atoms were located in subsequent Fourier syntheses.

The refinement was carried out by full-matrix leastsquares calculations with the program RFINE4 (Finger \& Prince, 1975). The quantity minimized in the leastsquares calculations was $\Sigma w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2}$, where $w=$ $\left[\sigma\left(F_{\text {ave }}\right)\right]^{-2} . \sigma^{2}\left(F_{\text {ave }}\right)$ is taken as $\Sigma_{i=1}^{n} \sigma_{c}^{2}\left(F_{i}\right) / n$ if it is greater than $\sum_{i-1}^{n}\left(F_{i}-F_{\text {ave }}\right)^{2} /(n-1)$, where $n$ is the number of equivalent reflections; otherwise it is the average of the two estimates. Scattering factors for all atoms were those for the neutral atoms taken from International Tables for X-ray Crystallography (1974). Two cycles of least-squares calculations with individual isotropic temperature factors led to an $R\left(=\Sigma \| F_{o} \mid-\right.$ $\left.\left|F_{c} \| / \Sigma\right| F_{o} \mid\right)$ of 0.068 and an $R_{\text {h }}\left\{=\mid \Sigma w\left(F_{o}-\right.\right.$ $\left.\left.F_{c}\right)^{2} /\left.\Sigma w F_{o}{ }^{2}\right|^{1 / 2}\right\}$ of 0.044 . Two additional cycles of refinement with anisotropic thermal parameters for all atoms reduced $R$ to 0.047 and $R_{x}$ to 0.034 (the corresponding values for all 1985 reflections are 0.058 and 0.036 respectively). The refinement was terminated with an average shift/error in the last cycle of $0 \cdot 1$ and a maximum value of 0.47 . There were no abnormal correlation coefficients $1-0.40$ between $\beta_{13}$ of $\operatorname{Sn}(1)$ and $\beta_{13}$ of $\operatorname{Sn}(2)$ was the largest $]$. No allowance for ex-
tinction was found necessary. The final atomic parameters are listed in Table 1.*

## Results and discussion

The crystal structures of $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ and $\mathrm{SnSO}_{4}$ (Donaldson \& Puxley, 1972) provide an opportunity to see how the packing of $\mathrm{Sn}^{11}$ ions and $X \mathrm{O}_{4}$ ions varies with the cation-to-anion ratio. The overall molecular packing is illustrated in Fig. 1. The two fairly large open channels at $x=\frac{1}{2}$ that run parallel to the $b$ axis are formed by Sn atoms arranged in a helical fashion. Similar open channels have been noticed in other $\mathrm{Sn}^{\prime \prime}$ compounds, e.g. $\mathrm{NaSn}_{2} \mathrm{~F}_{5}$ (McDonald, Larson \& Cromer, 1964), $\quad \mathrm{Sn}_{3} \mathrm{~F}_{3} \mathrm{PO}_{4}$ (Berndt, 1972), $\mathrm{Sn}_{3} \mathrm{O}(\mathrm{OH})_{2} \mathrm{SO}_{4}$ (Davies, Donaldson, Laughlin, Howie \& Beddoes, 1975), and $\mathrm{Sn}_{2} \mathrm{OHPO}_{4}$ (Jordan, Schroeder, Dickens \& Brown, 1976). Such open spaces are generally con-

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Fig. 1. A stereoscopic view of the unit cell of $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$. The origin of the coordinate system is marked by an asterisk.
sidered to be a result of the hemispherical coordination preferred by $\mathrm{Sn}^{11}$ ions. As in the case of $\mathrm{SnSO}_{4}$, the basic coordination of $\mathrm{Sn}^{11}$ can be clearly ascertained from the observation (see Table 2) that three $\mathrm{Sn}-\mathrm{O}$ distances are much shorter than the others. This eliminates the need for choosing a cut-off distance as a means of distinguishing between bonds of strong and intermediate strength.

The details of the coordination about $\operatorname{Sn}(1)$ and $\mathrm{Sn}(2)$ are shown in Fig. 2(a) and (b). $\mathrm{Sn}(1)$ is strongly coordinated to three O atoms from three different $\mathrm{P}(1) \mathrm{O}_{4}$ groups (average $\mathrm{Sn}-\mathrm{O}$ distance $2 \cdot 119 \AA$ ). The Sn atom is at the apex of a trigonal pyramid with the three O atoms at the base. The coordination around $\mathrm{Sn}(1)$ is completed by two additional weak contacts (average $\mathrm{Sn}-\mathrm{O}$ distance $2.971 \AA$ ).

The primary coordination around $\mathrm{Sn}(2)$ is quite similar to that of $\mathrm{Sn}(1)$. There are three strong $\mathrm{Sn} \cdots \mathrm{O}$ bonds (average distance $2 \cdot 100 \AA$ ) from three different $\mathrm{P}(2) \mathrm{O}_{4}$ groups. Again, Sn and the bonding O atoms form a trigonal pyramid. The only other $\mathrm{Sn} \cdots \mathrm{O}$ contact, $\mathrm{Sn}(2)-\mathrm{O}(22)$, is at a distance of $3.054 \AA$.

The $\mathrm{Sn}-\mathrm{O}$ distances and $\mathrm{O}-\mathrm{Sn}-\mathrm{O}$ angles (see Table 2) for $\operatorname{Sn}(1)$ and $\operatorname{Sn}(2)$ that make up the primary coordination geometry are quite similar to those observed in

Table 2. Bond angles $\left({ }^{\circ}\right)$ in $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$
Sn environment

| $\mathrm{O}(11)-\mathrm{Sn}(1)-\mathrm{O}(13)$ | $89.8(2)$ |
| :--- | :--- |
| $\mathrm{O}(11)-\mathrm{Sn}(1)-\mathrm{O}(14)$ | $88 \cdot 2(2)$ |
| $\mathrm{O}(13)-\mathrm{Sn}(1)-\mathrm{O}(14)$ | $86.8(2)$ |
| $\mathrm{O}(22)-\mathrm{Sn}(2)-\mathrm{O}(23)$ | $88.5(3)$ |
| $\mathrm{O}(22)-\mathrm{Sn}(2)-\mathrm{O}(24)$ | $88.3(3)$ |
| $\mathrm{O}(23)-\mathrm{Sn}(2)-\mathrm{O}(24)$ | $94.2(3)$ |
| $\mathrm{O}(21)-\mathrm{Sn}(3)-\mathrm{O}(12)$ | $84 \cdot 9(2)$ |
| $\mathrm{O}(21)-\mathrm{Sn}(3)-\mathrm{O}\left(2^{\prime}\right)$ | $84.8(2)$ |
| $\mathrm{O}(12)-\mathrm{Sn}(3)-\mathrm{O}\left(12^{\prime}\right)$ | $69.6(2)$ |

$\mathrm{PO}_{4}$ groups and environment
$\mathrm{O}(11)-\mathrm{P}(1)-\mathrm{O}(12)$
$\mathrm{O}(11)-\mathrm{P}(1)-\mathrm{O}(13)$
$\mathrm{O}(11)-\mathrm{P}(1)-\mathrm{O}(14)$
$\mathrm{O}(12)-\mathrm{P}(1)-\mathrm{O}(13)$
$\mathrm{O}(12)-\mathrm{P}(1)-\mathrm{O}(14)$
$\mathrm{O}(13)-\mathrm{P}(1)-\mathrm{O}(14)$
$\mathrm{O}(21)-\mathrm{P}(2)-\mathrm{O}(22)$
$\mathrm{O}(21)-\mathrm{P}(2)-\mathrm{O}(23)$
$\mathrm{O}(21)-\mathrm{P}(2)-\mathrm{O}(24)$
$\mathrm{O}(22)-\mathrm{P}(2)-\mathrm{O}(23)$
$\mathrm{O}(22)-\mathrm{P}(2)-\mathrm{O}(24)$
$\mathrm{O}(23)-\mathrm{P}(2)-\mathrm{O}(24)$
$\mathrm{P}(1)-\mathrm{O}(11)-\mathrm{Sn}(1)$
$\mathrm{P}(1)-\mathrm{O}(12)-\mathrm{Sn}(3)$
$\mathrm{P}(1)-\mathrm{O}(12)-\mathrm{Sn}\left(3^{\prime}\right)$
$\mathrm{P}(1)-\mathrm{O}(13)-\mathrm{Sn}(1)$
$\mathrm{P}(1)-\mathrm{O}(14)-\mathrm{Sn}(1)$
$\mathrm{P}(2)-\mathrm{O}(21)-\mathrm{Sn}(3)$
$\mathrm{P}(2)-\mathrm{O}(22)-\mathrm{Sn}(2)$
$\mathrm{P}(2)-\mathrm{O}(23)-\mathrm{Sn}(2)$
$\mathrm{P}(2)-\mathrm{O}(24)-\mathrm{Sn}(2)$
110.8 (4)
113.6 (4)
107.3 (4)
103.7 (4)
110.0 (4)
111.5 (4)
$1110(4)$
108.5 (4)
107.7 (4)
$113.0(4)$
106.5 (4)
110.0 (4)
135.9 (4)
135.6 (4)
110.7 (4)
141.0 (4)
117.2 (4)
121.5 (4)
$135.9(4)$
139.9 (4)
119.4 (4)
a wide variety of $\mathrm{Sn}^{11}$ compounds (for surveys see Jordan, Schroeder, Dickens \& Brown, 1976; McDonald, Hau \& Eriks, 1976).

The $\mathrm{O}-\mathrm{Sn}-\mathrm{O}$ angles average $10^{\circ}$ larger than the corresponding ones in $\mathrm{SnSO}_{4}$. In $\mathrm{SnSO}_{4}$. nine weakly bonded O atoms surround each Sn in addition to the three strongly bonded O atoms. $\operatorname{In} \mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$, the number of weakly bonded O atoms to $\mathrm{Sn}(2)$ and $\mathrm{Sn}(1)$ are one and two respectively. This suggests that the additional $\mathrm{O} \cdots \mathrm{O}$ repulsions in $\mathrm{SnSO}_{4}$ may tend to compress the primary coordination geometry of $\mathrm{Sn}^{11}$.


Fig. 2. (a) The environment of $\mathrm{Sn}(1)$ in $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$. (b) The environment of $\mathrm{Sn}(2)$. Standard deviations for the distances given are $0.007 \AA$.

The coordination around $\mathrm{Sn}(3)$ is different from those of the other two Sn atoms. The trigonal coordination is maintained, $\bar{d}(\mathrm{Sn}-\mathrm{O})=2.231 \AA$, but there are three weaker bonds, $\bar{d}(\mathrm{Sn}-\mathrm{O})=2.888 \AA$, and one very weak bond ( $3.250 \AA$ ). The salient feature of this coordination is the oxygen-bridged $\operatorname{Sn}(3) \cdots \operatorname{Sn}(3)$ dimer lying across a center of symmetry (Fig. 3). The unusually small angle, $\mathrm{O}(12)-\mathrm{Sn}(3)-\mathrm{O}\left(12^{\prime}\right)=$ $69 \cdot 6(2)^{\circ}$, may be a result of the bridge formation. The small $\mathrm{O}-\mathrm{Sn}-\mathrm{O}$ angle and the $\mathrm{Sn} \cdots \mathrm{O}$ bond distances of 2.230 and $2.317 \AA$, the longest ones in the primary coordination of $\mathrm{Sn}^{\prime \prime}$ in $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$, indicate that $\mathrm{Sn} \cdots \mathrm{Sn}$ repulsion plays an important part in determining the dimeric configuration observed here. Such dimers have also been found in $\mathrm{NaSn}_{2} \mathrm{~F}_{5}$ (McDonald, Larson \& Cromer, 1964). As is the case with the other Sn atoms in this structure, the hemispherical coordination is maintained; the open side near the channel formed by $\mathrm{Sn}(1)$ and $\mathrm{Sn}(2)$ is presumably the site of the lone pair of electrons on $\operatorname{Sn}(3)$.

Fig. 1 also shows that the $\mathrm{Sn}^{11}$ and $\mathrm{PO}_{4}$ ions form alternating layers parallel to the $b c$ plane. The $\mathrm{P}(1) \mathrm{O}_{4}$ group layers can be thought of as made up of zigzag chains of $\mathrm{PO}_{4}$ groups parallel to the $c$ axis. Layers of $\mathrm{P}(1) \mathrm{O}_{4}$ groups are linked to layers of $\mathrm{P}(2) \mathrm{O}_{4}$ ions by $\mathrm{Sn}(3)$. Planes parallel to (100) might show cleavage because only one strong bond, $\operatorname{Sn}(3) \cdots \mathbf{O}(21)$, needs to be broken in separating the layers.

Although $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ is a product of the reaction of stannous ions with hydroxyapatite, its crystal structure shows no significant relationship to that of


Fig. 3. The environment of $\operatorname{Sn}(3)$. The $\mathrm{O}(12) \cdots \mathrm{O}(12)$ edge of the trigonal pyramid is shared with $\operatorname{Sn}\left(3^{\prime}\right)$, which is related to $\operatorname{Sn}(3)$ by a center of symmetry. Standard deviations for the distances given are $0.007 \AA$.
$\mathrm{Ca}_{5}\left(\mathrm{PO}_{4}\right)_{3} \mathrm{OH}$ (Kay, Young \& Posner, 1964). This suggests that $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ is unlikely to be formed as a crystalline overgrowth on hydroxyapatite upon reaction with stannous ions. In fact, the structure of $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ seems more related to those calcium phosphates that have layers of $\mathrm{PO}_{4}$ groups alternating with calcium ions and water molecules, e.g. $\mathrm{Ca}\left(\mathrm{H}_{2} \mathrm{PO}_{4}\right)_{2} \cdot \mathrm{H}_{2} \mathrm{O}$ (Schroeder, Prince \& Dickens, 1975).

The geometry of the $\mathrm{PO}_{4}$ groups in $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{2}$ can be interpreted in terms of the ionic packing and the coordination of the O atoms. The $\mathrm{P}-\mathrm{O}$ distances vary from 1.510 (7) to 1.544 (7) $\AA$ with a mean value of $1.535 \AA$, compared to $1.536 \AA$ calculated for 64 orthophosphate groups (Baur, 1974). All O atoms except $\mathrm{O}(12)$ are strongly bonded to one Sn atom; $\mathrm{O}(12)$ forms two strong bonds via bridging. None of these primary contacts involve edge coordination of a $\mathrm{PO}_{4}$ group. Each O atom except $\mathrm{O}(12)$ is also involved in a second weak $\mathrm{Sn} \cdots \mathrm{O}$ bond (Figs. 2 and 3). When weak $\mathrm{Sn} \cdots \mathrm{O}$ coordination is also taken into account, there are three $\mathrm{PO}_{4}$ shared-edge coordinations, one for each Sn atom. In each case the associated $\mathrm{O}-\mathrm{P}-\mathrm{O}$ angle is significantly less than the ideal tetrahedral angle. The primary contact $\mathrm{Sn}(3)-\mathrm{O}(12)$ at $2.317 \AA$ and the secondary contact $\mathrm{Sn}(3)-\mathrm{O}(13)$ at $2 \cdot 856 \AA$ result from edge coordination of $\mathrm{P}(1) \mathrm{O}_{4}$. The lengths of these contacts are more nearly equal than in the other cases of $\mathrm{PO}_{4}$ edge coordination and the $\mathrm{O}(12)-\mathrm{P}(1)-\mathrm{O}(13)$ angle is the smallest. This shows that the extent of the $\mathrm{PO}_{4}$ group distortions is also dependent on the secondary coordination of the O atoms.

The $\mathrm{P}-\mathrm{O}-\mathrm{Sn}$ angles for the primary coordinations vary widely over the range $111-141^{\circ}$ (see Table 2) with a mean value of $128 \cdot 5^{\circ}$. The four $\mathrm{P}-\mathrm{O}-\mathrm{Sn}$ angles in the range $111-122^{\circ}$ might be taken as indicating a tendency toward $s p^{2}$ hybridization at the O atom. However, there are also four $\mathrm{P}-\mathrm{O}-\mathrm{Sn}$ angles in the narrower range $135-141^{\circ}$. Values in the neighborhood of $120^{\circ}$ would also be expected from the repulsion of the P and two Sn atoms coordinated to the O . The wide variation of the $\mathrm{P}-\mathrm{O}-\mathrm{Sn}$ angles suggests that maintaining a strict bond direction at the O atom is not an important factor in determining the ionic arrangements in $\mathrm{Sn}_{3}\left(\mathrm{PO}_{4}\right)_{3}$.

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# The Structure of the Crystalline Complex of Purine and Urea (2:1) 

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#### Abstract

$\mathrm{C}_{5} \mathrm{H}_{4} \mathrm{~N}_{4} \cdot \frac{1}{2}\left(\mathrm{CH}_{4} \mathrm{~N}_{2} \mathrm{O}\right.$ ), FW 150.14, is orthorhombic, space group Fdd2, $a=20.921$ (4), $b=35 \cdot 274$ (5), $c=3.622$ (1) $\AA, V=2672.9 \AA^{3}, Z=16, D_{m}=1.490, D_{x}=1.496 \mathrm{~g} \mathrm{~cm}^{-3}$. The purine molecules are stacked along the $c$ axis and are held together by $\mathrm{N}(9)-\mathrm{H} \cdots \mathrm{N}(1)$ and N (urea) $-\mathrm{H} \cdots \mathrm{N}(7)$ hydrogen bonds. The protonation site in the imidazole ring is $N(9)$ instead of the usual $N(7)$ in purine. The overlapping mode of the purine molecules differs from that in the crystal of purine. The urea molecules are also associated with each other through hydrogen bonding. Two half-weight urea molecules are arranged around the diad axis parallel to [001] with their O atoms on the diad and their molecular planes nearly parallel to ( 010 ). Because of the short axial length of $c$, the repeat distance of the urea molecules must be a multiple of $c$, thus indicating disorder along the $c$ axis.


## Introduction

We found that purine and urea form a stable 2:1 molecular complex in the solid state. The complex retains the same stoichiometric ratio through sublimation at $105^{\circ} \mathrm{C} / 23 \mathrm{mmHg}$ as well as recrystallization from various organic solvents. The infrared spectrum (in KBr ) of the complex shows a drastic change in the $1550 \sim 1700 \mathrm{~cm}^{-1}$ region compared with the mixture of the same composition. At present, there is no evidence of the analogous complex formation in solution. In order to clarify the state of association between purine and urea in this molecular complex, we have carried out the crystal structure analysis.

## Experimental

Crystals of the complex were prepared by dissolving purine and urea in a $2: 1$ stoichiometric ratio in warm ethyl acetate and allowing the mixture to stand at room temperature. A crystal $0.3 \times 0.1 \times 0.45 \mathrm{~mm}$ was cut with a razor and used for the X-ray experiments. X-ray data were collected on a Philips PW 1100 automatic diffractometer using Cu Ka radiation monochromated by a graphite plate. The systematic absences (Table 1) established the orthorhombic, face-centred space group $F d d 2$. The calculated density, assuming 16 structural units in a cell, accorded well with that measured in $\mathrm{CCl}_{4}-n$-hexane by flotation. This indicates that the urea


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[^1]:    * A list of structure factors has been deposited with the British Library Lending Division as Supplementary Publication No. SUP 32340 ( 8 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 13 White Friars, Chester CHI 1NZ, England.

